

*It's all about GARY*



*how Physicists  
celebrate a birthday*

*John Dabson*

# CONTENTS

*Memories: Gary gives Permission  
to Be Creative*

*Something New: Event Chirality,  
or "Handicity", and How to Sort It*

*Strange Features of the  
Two Dimensional World*

*Analytic Continuation:  
What Fun is THAT!*

*And then, we Celebrate !*

#### 14) Exchange Degeneracy in K<sup>+</sup> Lambda Or Sigma<sub>0</sub> Photoproduction.

Gary R. Goldstein, (Tufts U.) . Print-74-0854 (TUFTS), Mar 1974. 5pp.

Published in **Nucl.Phys.B79:341,1974.**

[References](#) | [LaTeX\(US\)](#) | [LaTeX\(EU\)](#) | [Harvmac](#) | [BibTeX](#) | Cited 1 time

Journal Server [doi:[10.1016/0550-3213\(74\)90491-X](https://doi.org/10.1016/0550-3213(74)90491-X)]

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#### 15) Optimally Simple Connection Between the Reaction Matrix and the Observables.

Gary R. Goldstein, (Tufts U.) , Michael J. Moravcsik, (Oregon U.) . PRINT-74-0982 (OREGON), May 1974. (Published 54pp.

Published in **Annals Phys.98:128,1976.**

TOPCITE = 50+

[References](#) | [LaTeX\(US\)](#) | [LaTeX\(EU\)](#) | [Harvmac](#) | [BibTeX](#) | [Keywords](#) | Cited 62 times

Journal Server [doi:[10.1016/0003-4916\(76\)90241-4](https://doi.org/10.1016/0003-4916(76)90241-4)]

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#### 16) Semilocal Duality in pi0 Photoproduction.

H.K. Armenian, Gary R. Goldstein, J.P. Rutherford, D.L. Weaver, (Tufts U.) . PRINT-74-1685 (TUFTS), Sep 1974. 15p

Published in **Phys.Rev.D12:1278,1975.**

[References](#) | [LaTeX\(US\)](#) | [LaTeX\(EU\)](#) | [Harvmac](#) | [BibTeX](#) | [Keywords](#) | Cited 1 time

Journal Server [doi:[10.1103/PhysRevD.12.1278](https://doi.org/10.1103/PhysRevD.12.1278)]

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Gary and Mike:

Experiments that average over observable information,

like spin, don't really tell you much:

they don't test much, and they tend to confirm wrong theory



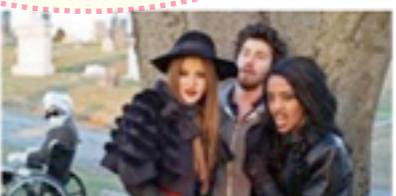
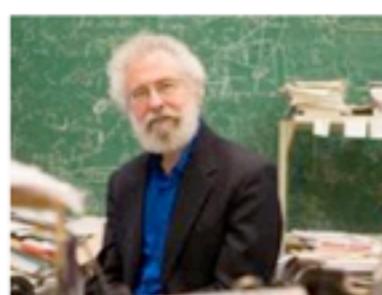
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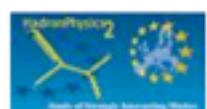
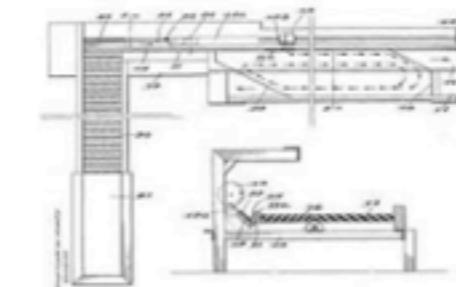
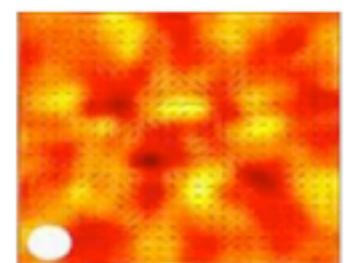
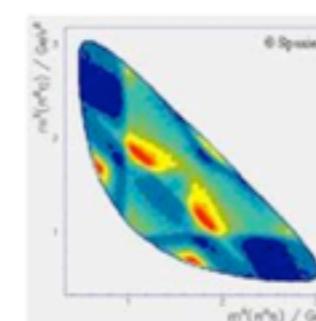
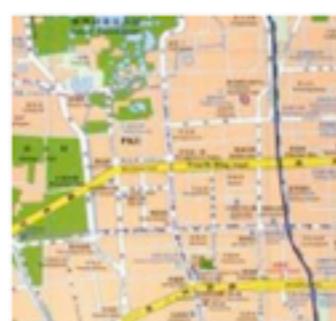
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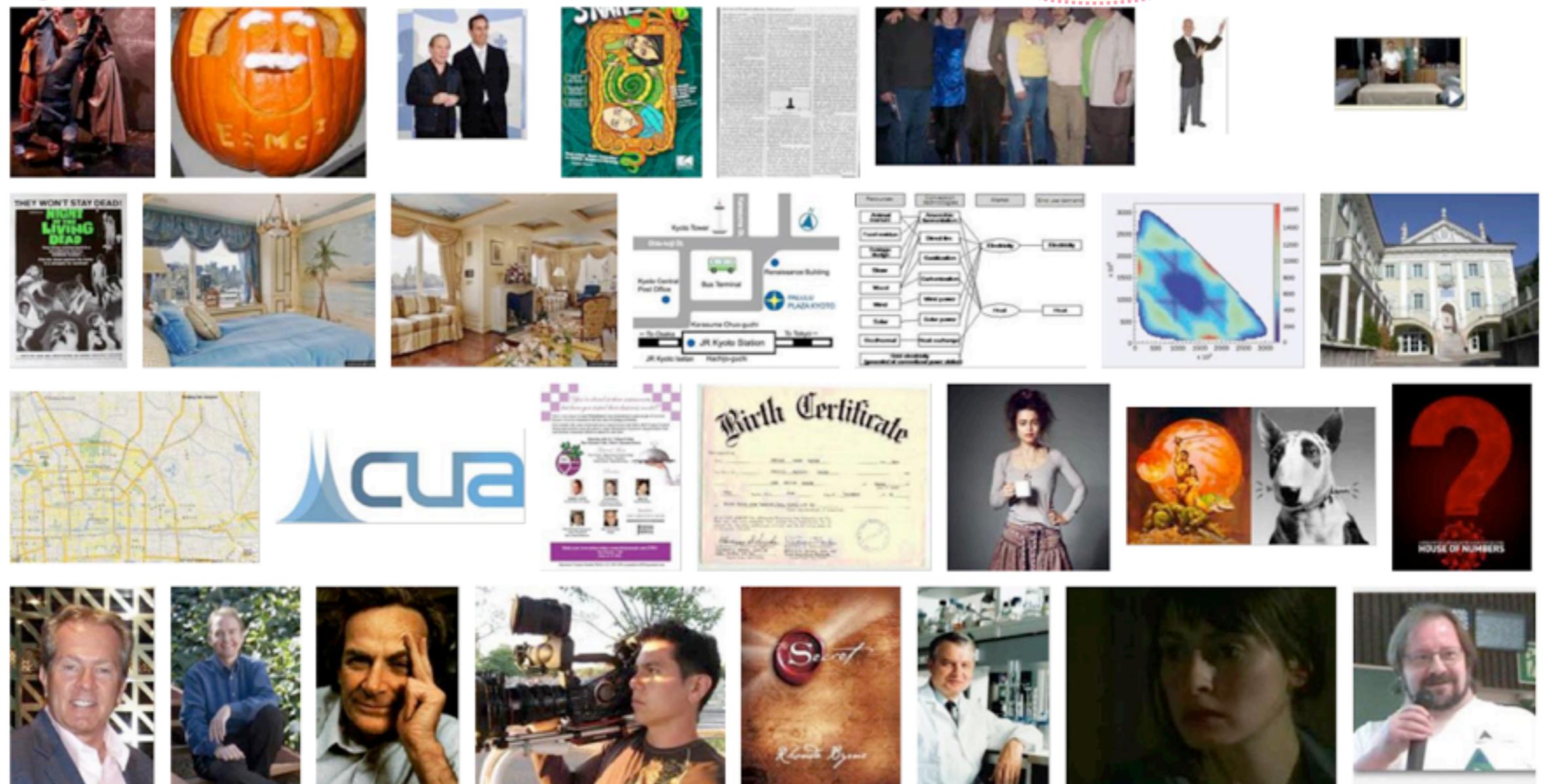
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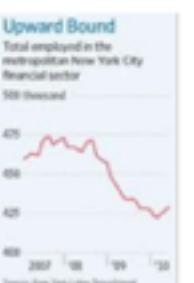
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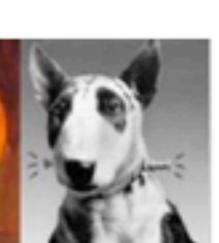
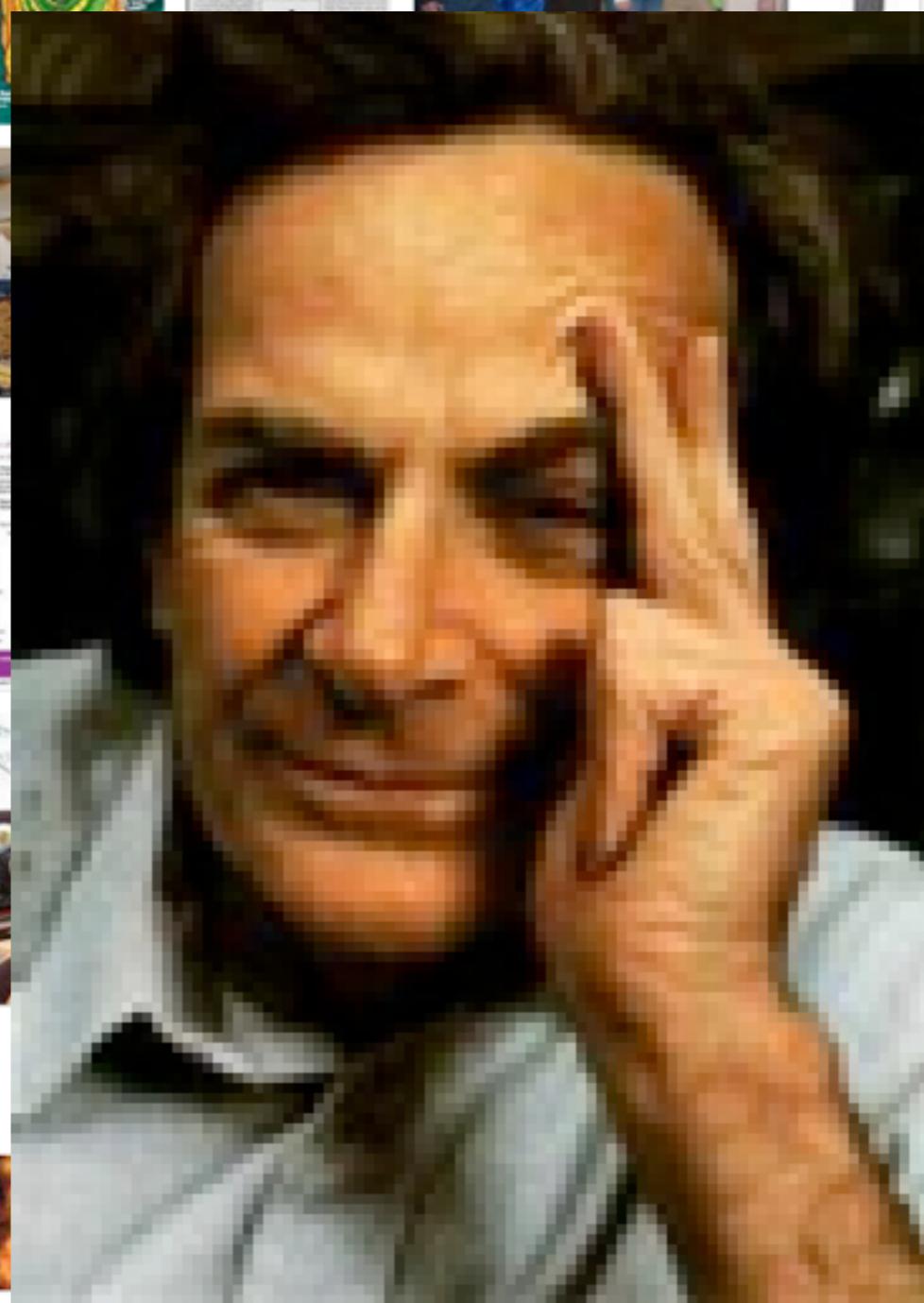
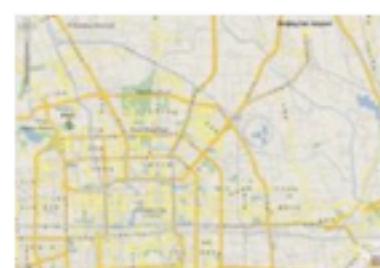
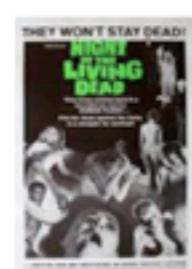
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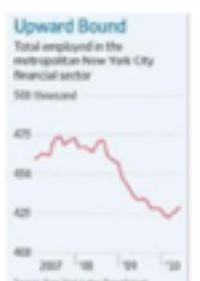
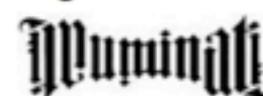
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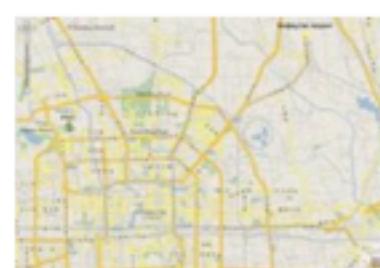
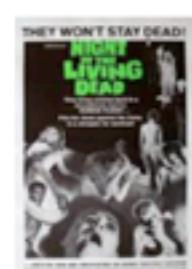
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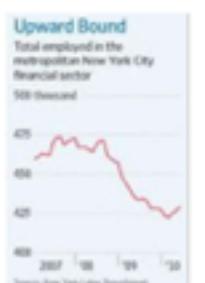
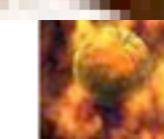
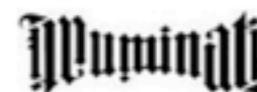
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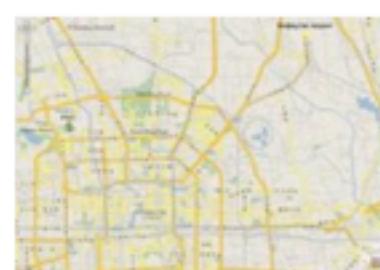
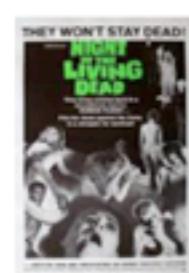
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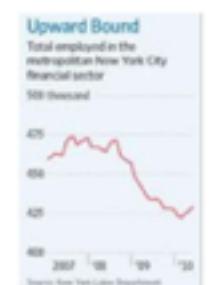
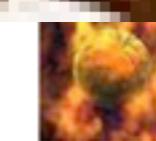
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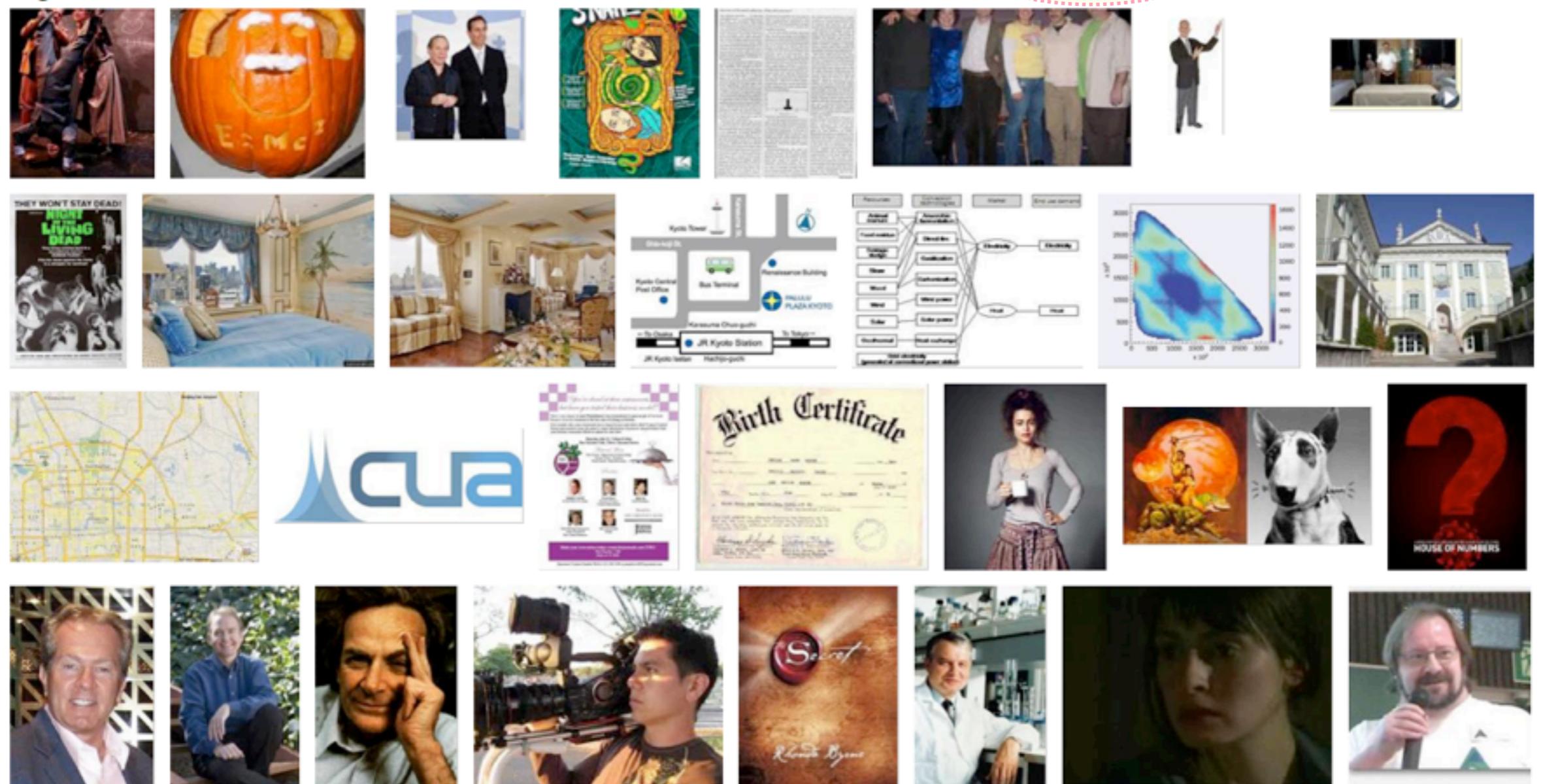
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# The Gary Effects:



we think a party is... Doing Physics



it's more than OK to  
think about SPIN



Always break azimuthal symmetry  
(there's so much to learn)



we think a party is ... Doing Physics....

Let's talk about...  
event-by-event  
parity, chirality,  
“handicity”



work with Mihailo  
BACKOVIC

Let's talk about event-by-event parity

$$D \in O(3); \quad \det D = -1$$

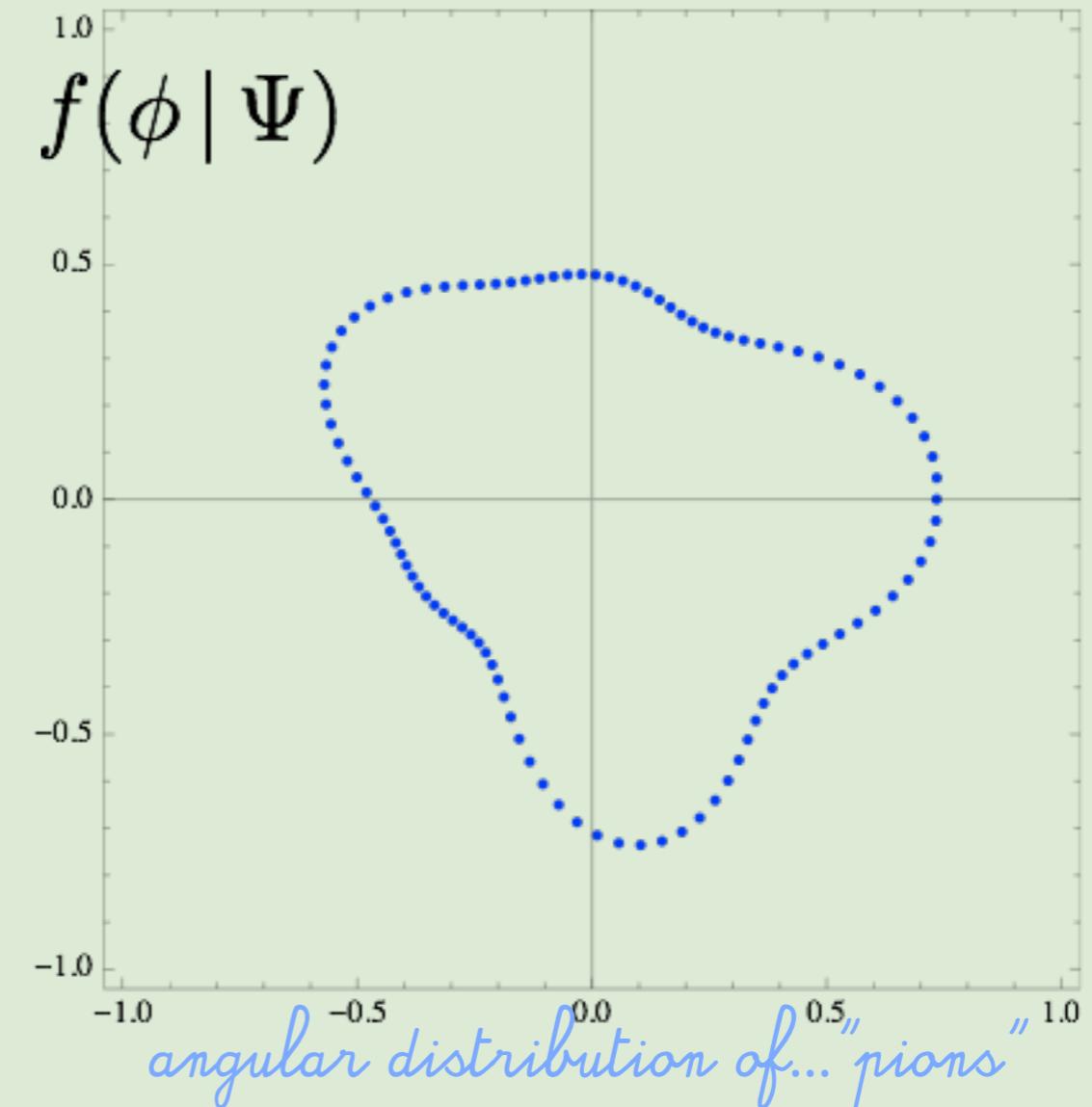
$$f(\phi) = \frac{dN}{d\phi}$$

$$f(\phi, \Psi) = \frac{dN}{d\phi d\Psi}$$

$$f(\phi, \Psi) = f(\phi | \Psi) f(\Psi)$$

for RHIC physics;  
Voloshin et al

Every event breaks parity symmetry



parity is NOT cosine/sine separation - that's origin-dependent

$$\frac{dN}{d\phi} = \frac{1}{2\pi} + v \cos\phi + a \sin\phi;$$

Let  $v = \rho \cos\Psi$ ;  $a = \rho \sin\Psi$ ;

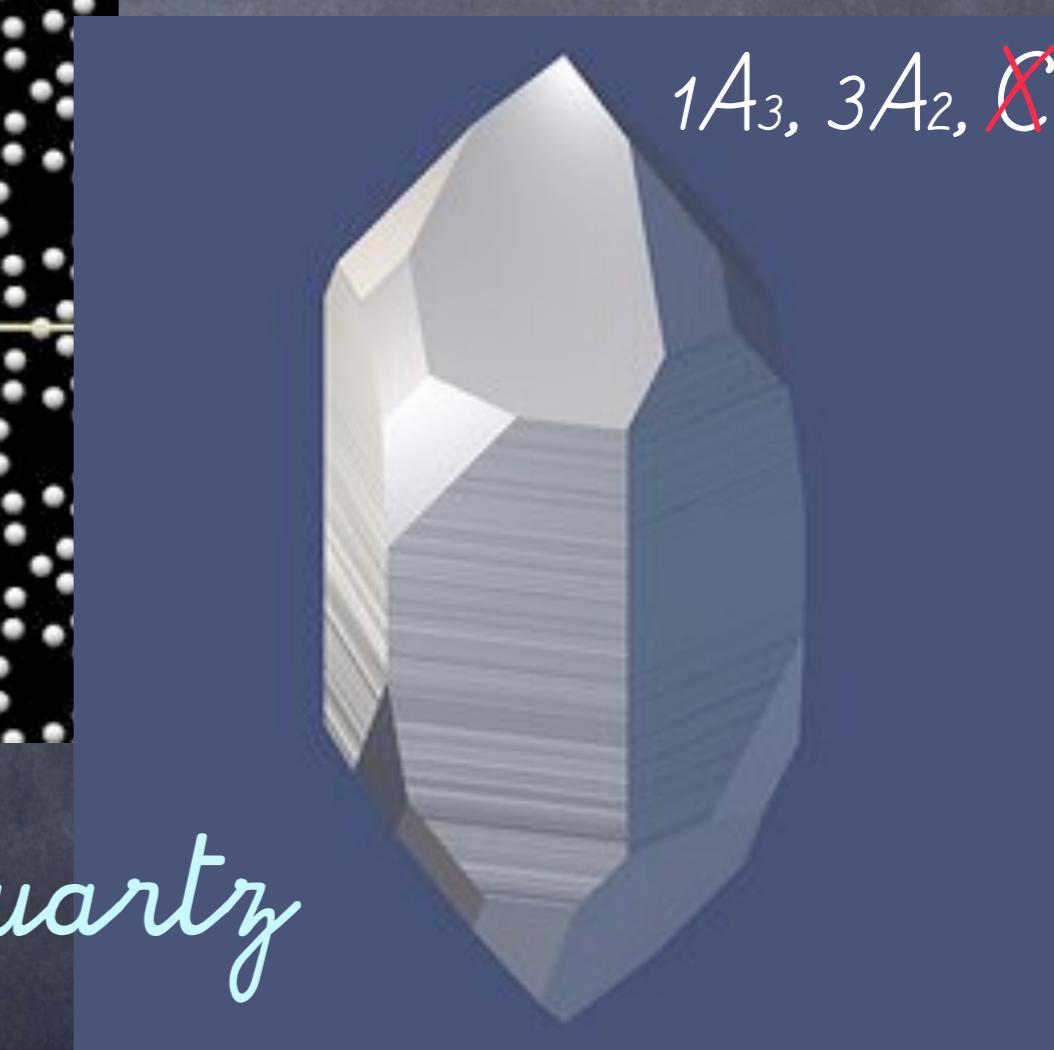
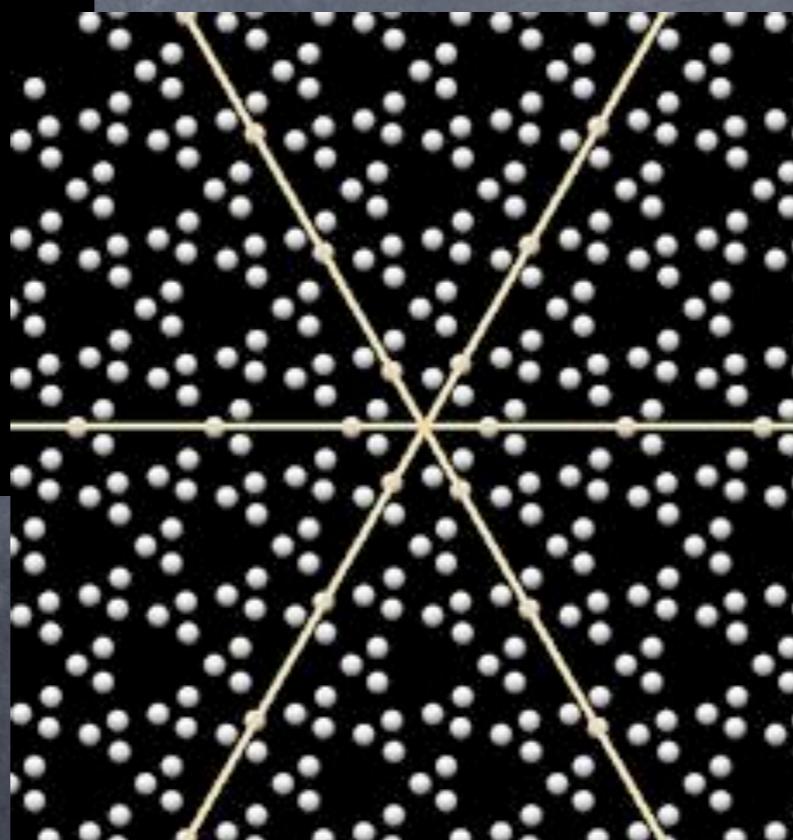
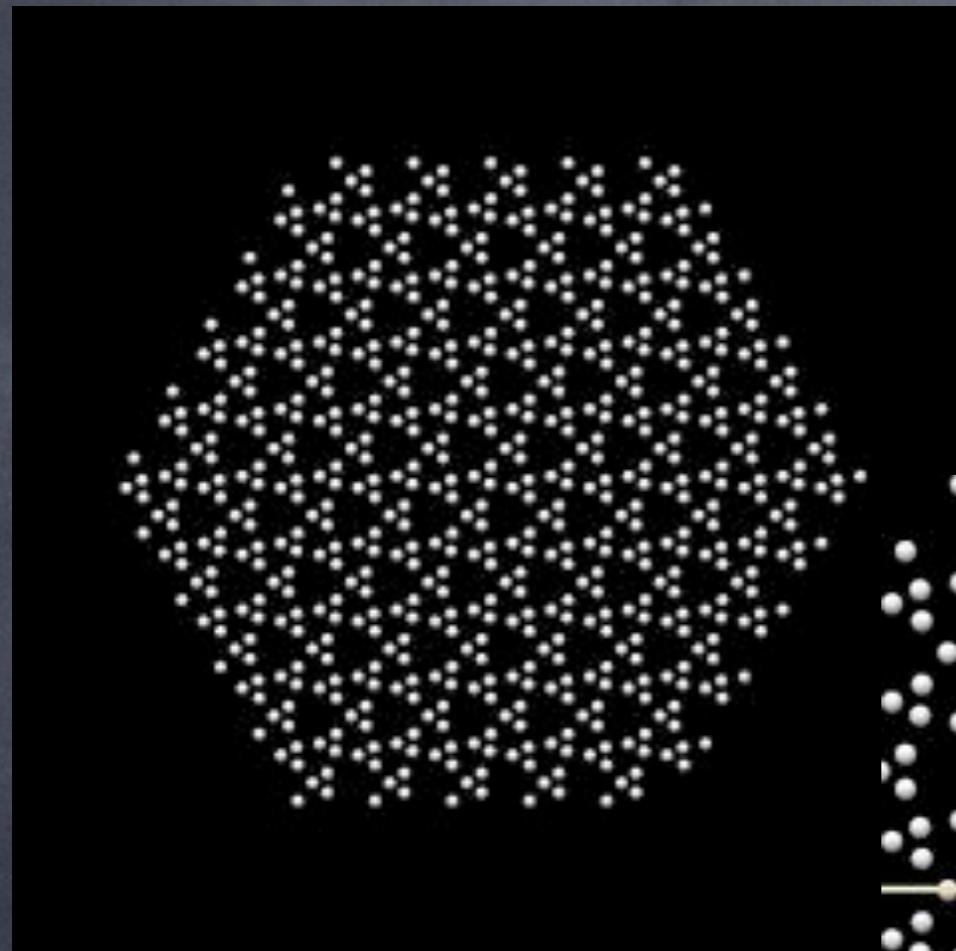
then  $\frac{dN}{d\phi} = \frac{1}{2\pi} + \rho \cos(\phi - \Psi)$

measures nothing about parity:  
STAR experiment

parity and event chirality have been overlooked...

spontaneous parity violation is COMMON in Nature,  
even macroscopically

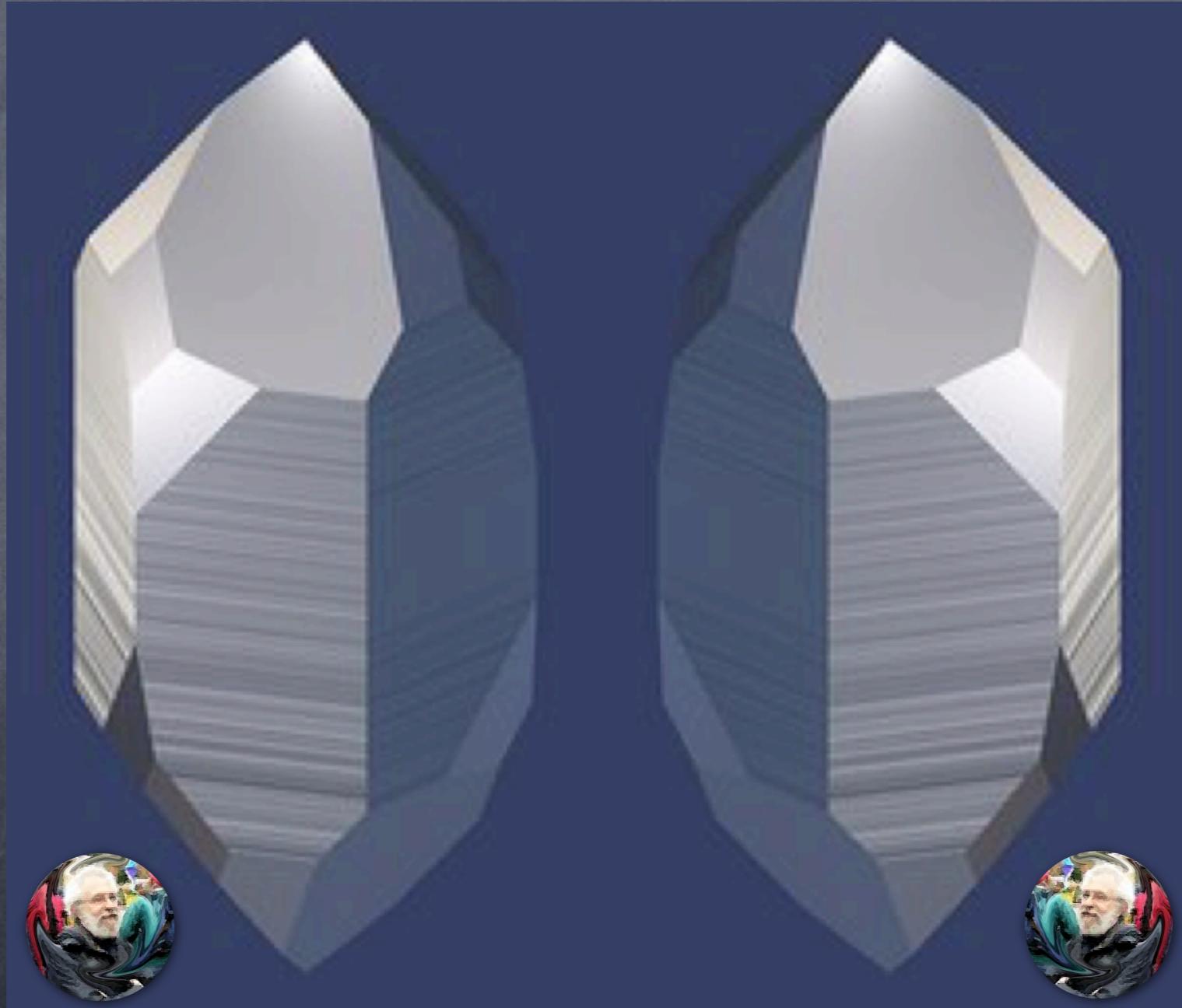
$$f(\text{object} \mid \text{chirality}) \neq 0$$



left and right handed quartz,  $\text{SiO}_2$

the unsorted,  
global, angle and  
chirality  
AVERAGED  
crystal is a...ball

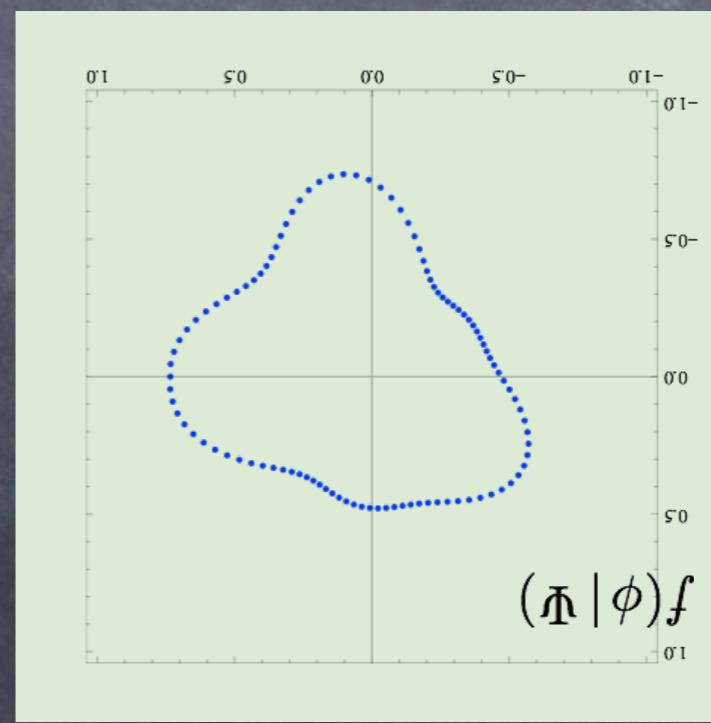
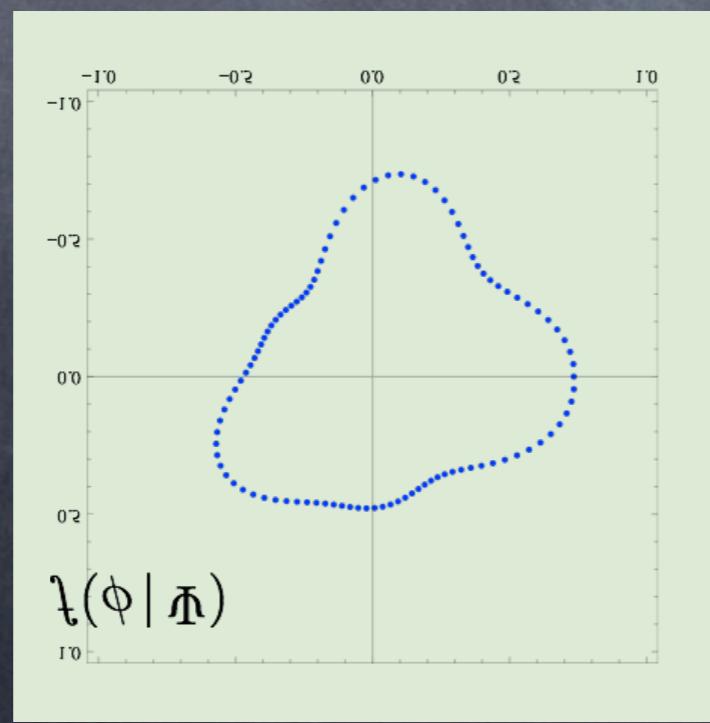
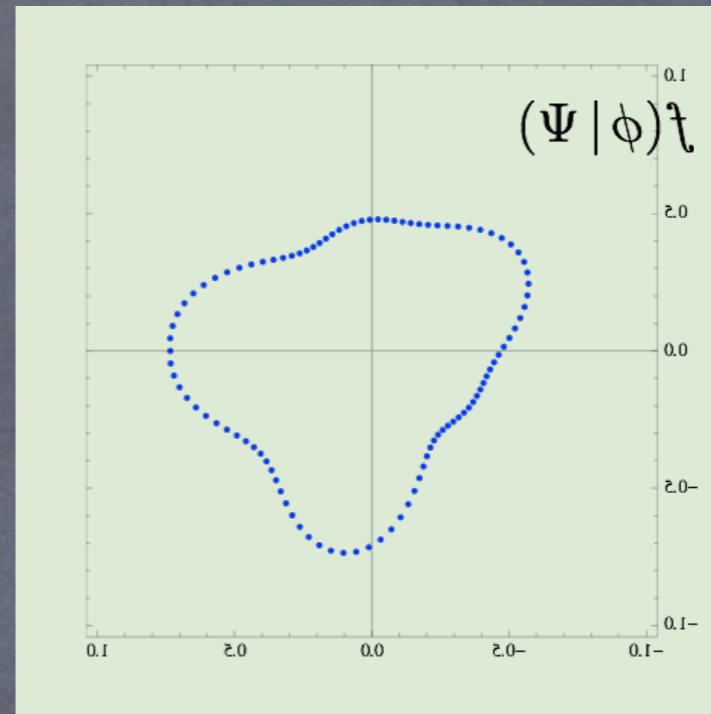
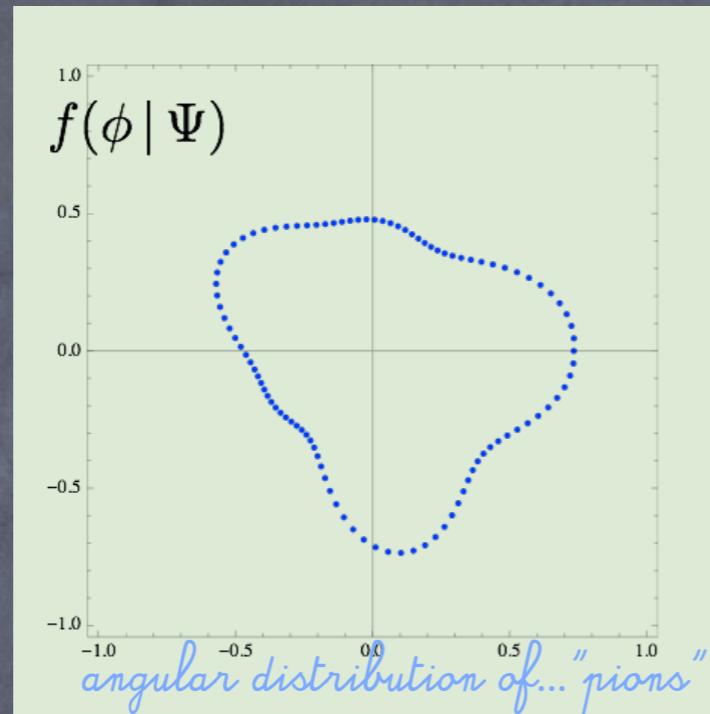
sort first;  
average  
afterwards



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“chirality sorting”: four events conditional on orientation



first sort,  
then add

events

# Violating 2D parity violates 3D parity

Consider any  $D \in O(3)$  with  $\det(D) = -1$ . If symmetry of the Hamiltonian  $H$  under  $D$  fails then  $[H, D] \neq 0$ . In three dimensions parity  $P = -1_{3 \times 3}$ . From  $D$  make a pure rotation  $R = P D$ ,  $\det(R) = 1$ . Then

$$[H, D] = [H, PR] = P[H, R] + [H, P]R \neq 0,$$

or

$$[H, P] \neq 0,$$

given rotational invariance  $[H, R] = 0$ . Then  $P$ -symmetry is violated by finding *any single case of parity-violation on a subspace*.

Yet 2D parity is far more complicated...

3 dimensions

$$\vec{x} \xrightarrow{P} -\vec{x}$$

$$\det \begin{pmatrix} -1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{pmatrix} = -1$$

2 dimensions

$$\vec{x} \xrightarrow{?} -\vec{x}$$

$$\det \begin{pmatrix} -1 & 0 \\ 0 & -1 \end{pmatrix} = +1$$

this is a rotation

There DOES NOT EXIST a rotationally invariant parity  
on TWO dimensional space

$D_N$ , the dihedral groups of index  $N$ ,  
form the discrete subgroups of  $O(2)$

generator algebra:

$$\mathcal{R}^N = 1, \quad \mathcal{P}^2 = 1, \quad \mathcal{P}\mathcal{R}\mathcal{P} = \mathcal{R}^{-1}$$

$2N$  elements  $(R_{Nk}, P_{Nk})$ , for  $k = (0\dots N-1)$ ,

with  $\det(R_{Nk}) = 1$ ,  $\det(P_{Nk}) = -1$

non-Abelian: no  $P_{Nk}$  commutes with any  $R_{Nk}$

all previous measures of 2D “chirality” in computer science, image  
processing...psychology...either fail under  $D_N$  or  $SO(2)$

$D_N$  elements:  $k=0 \dots N-1$

$$P_{Nk} = \begin{pmatrix} \cos 2\pi k/N & \sin 2\pi k/N \\ \sin 2\pi k/N & -\cos 2\pi k/N \end{pmatrix};$$

$$R_{Nk} = \begin{pmatrix} \cos 2\pi k/N & -\sin 2\pi k/N \\ \sin 2\pi k/N & \cos 2\pi k/N \end{pmatrix}.$$

how shall we measure and sort  
the degree of chirality in two dimensions?



the whirlyness of the classical distribution...

$$\int_0^{2\pi} d\phi \ f(-i\frac{\partial}{\partial\phi})f \quad \dots \text{equals zero.}$$

In quantum mechanics...

$$\int_0^{2\pi} d\phi \ \psi^*(-i\frac{\partial}{\partial\phi})\psi \neq 0$$

2D rotationally invariant;  $D_N$  (even, odd)



to get  $\psi$  from  $f$  ... analytically continue

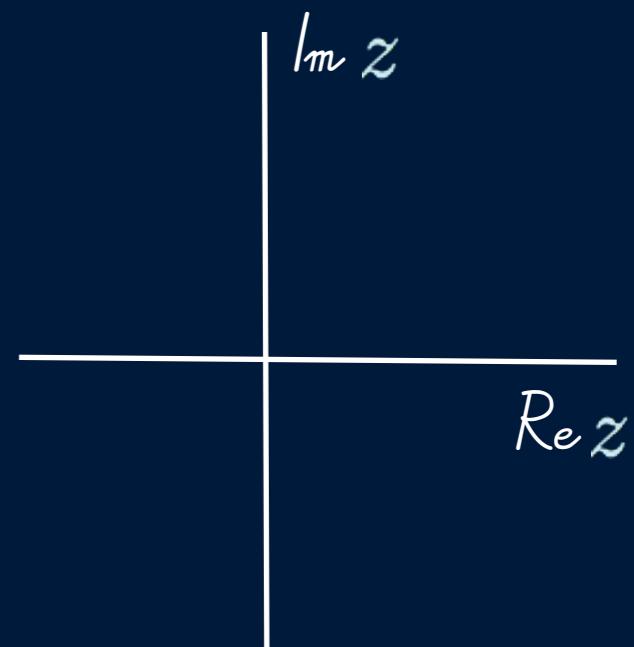
Analytic continuation:  $z = e^{i\phi}$

$$f(\phi) = \sum_{m=-\ell}^{\ell} f_m e^{im\phi} = z^{-\ell} \sum_{m=0}^{2\ell} f_m z^m,$$
$$= f_0 z^{-\ell} \prod_{k=1}^{2\ell} (z - z_k).$$

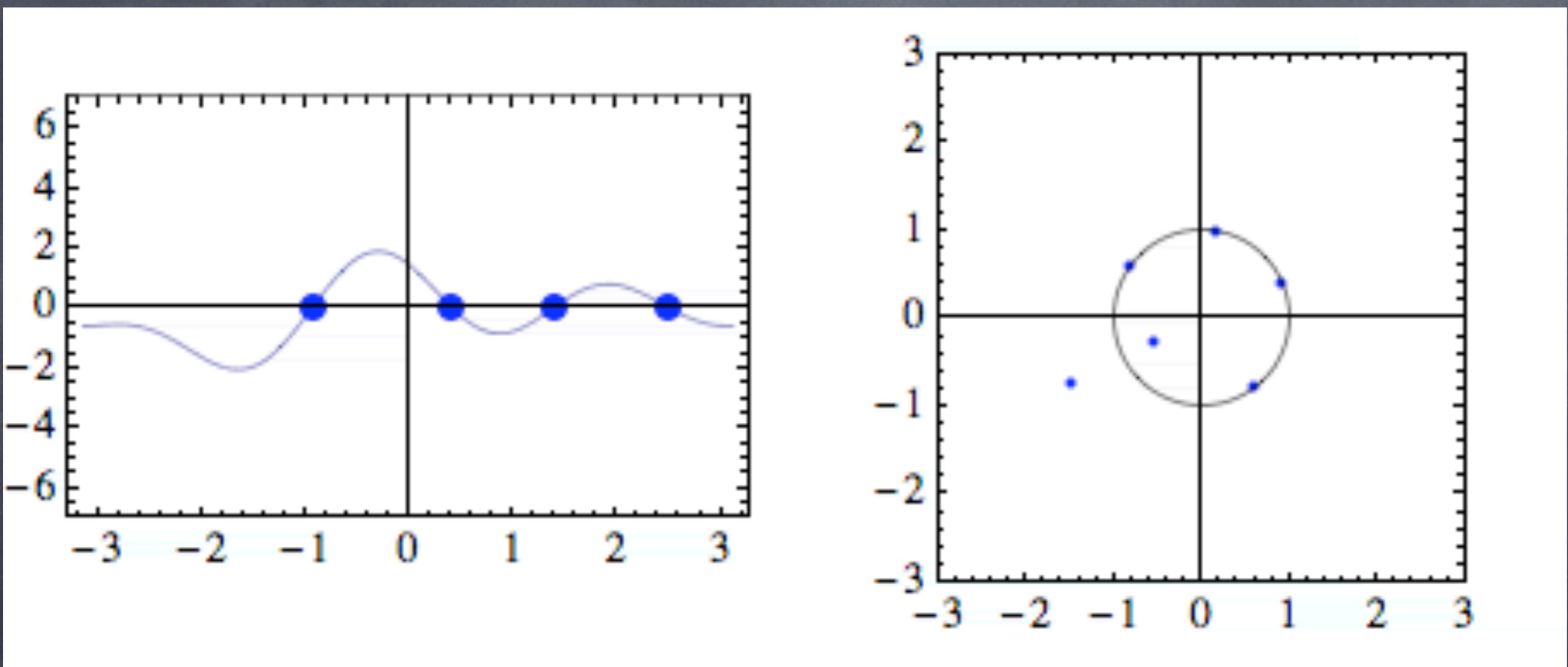
“fundamental  
theorem of  
algebra”



poles and zeroes



under rotations, complex zeros transform  
like 2-vectors



under rotations, complex zeros transform  
like 2-vectors

## Continuation, continued:

$$\begin{aligned} f(\phi) &= \sum_{m=-\ell}^{\ell} f_m e^{im\phi} = z^{-\ell} \sum_{m=0}^{2\ell} f_m z^m, \\ &= f_0 z^{-\ell} \prod_{k=1}^{2\ell} (z - z_k). \end{aligned}$$

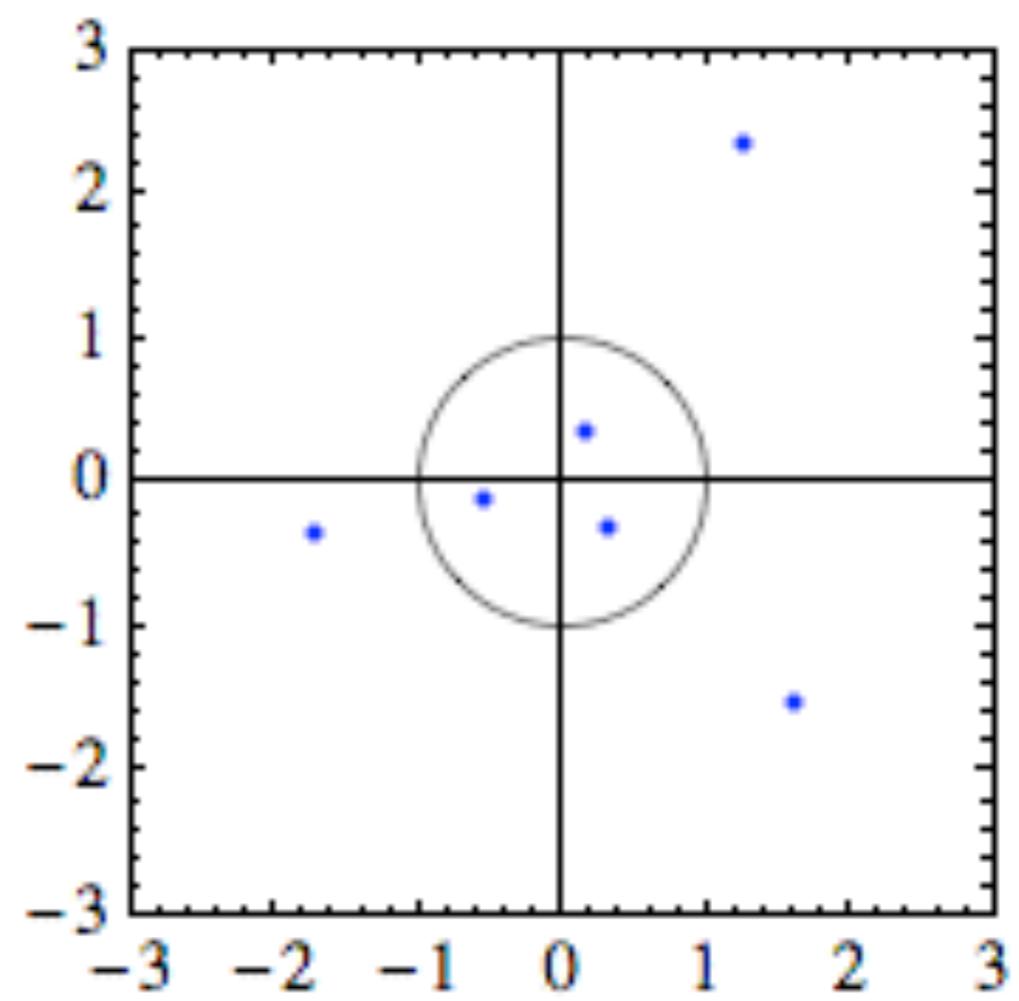
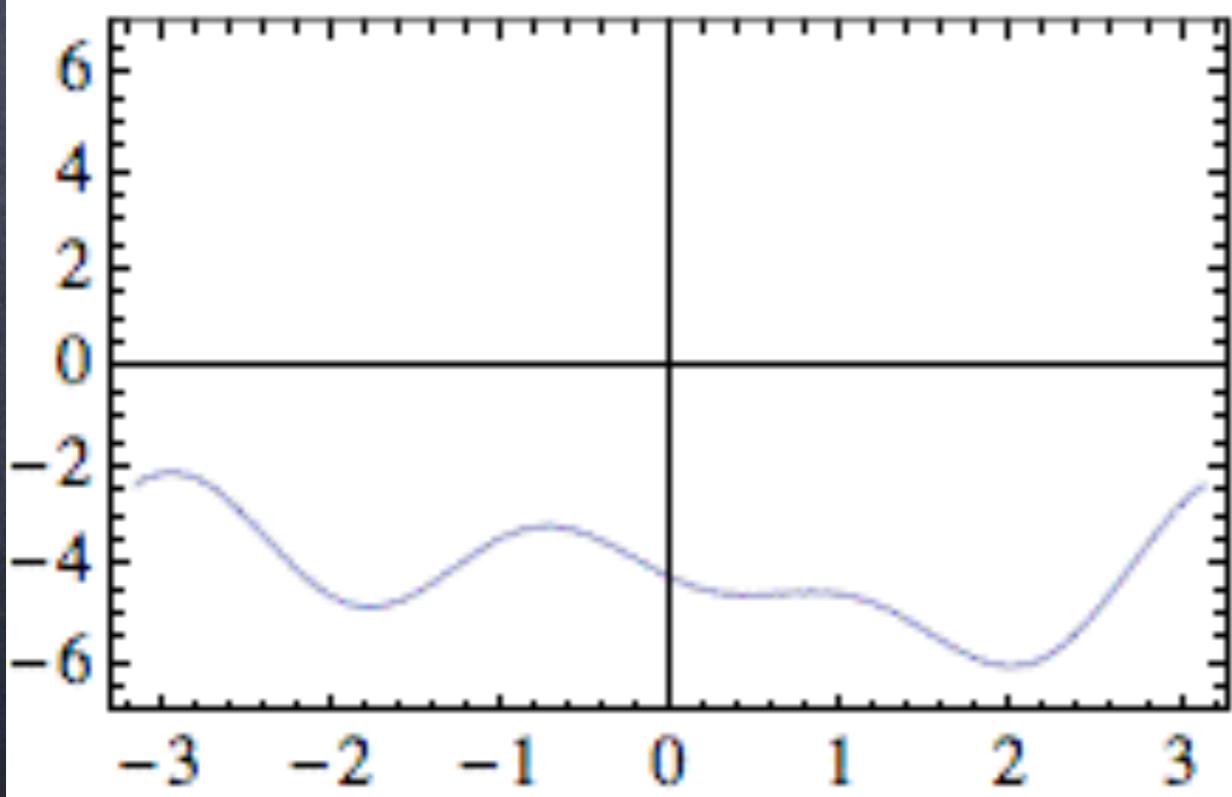
$f(\phi) = \text{real} \rightarrow \text{zeroes in pairs } (z_k, 1/z_k^*)$

$$\begin{aligned} f(z; |z| = 1) &= \psi(z)^* \psi(z) \\ \text{where } \psi(z) &= \psi_0 \prod_k \frac{z - z_k}{z} \end{aligned}$$

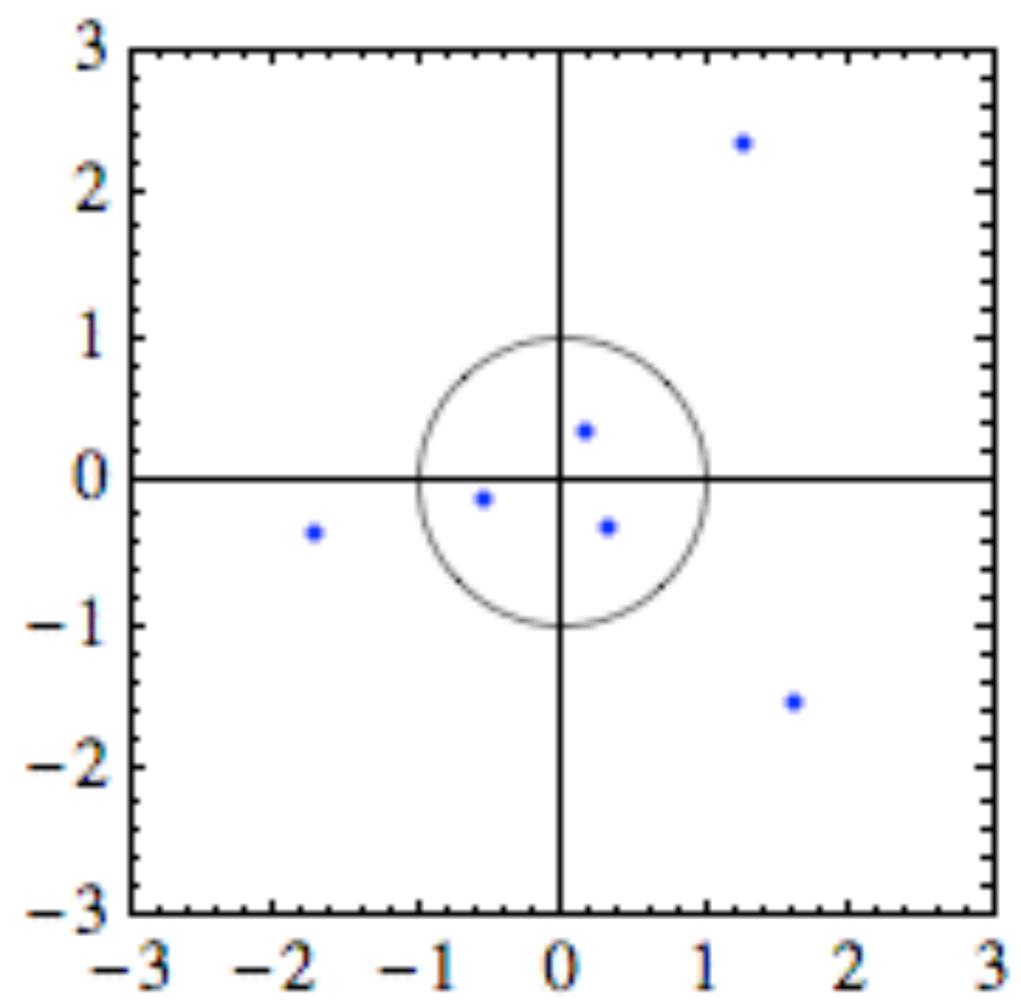
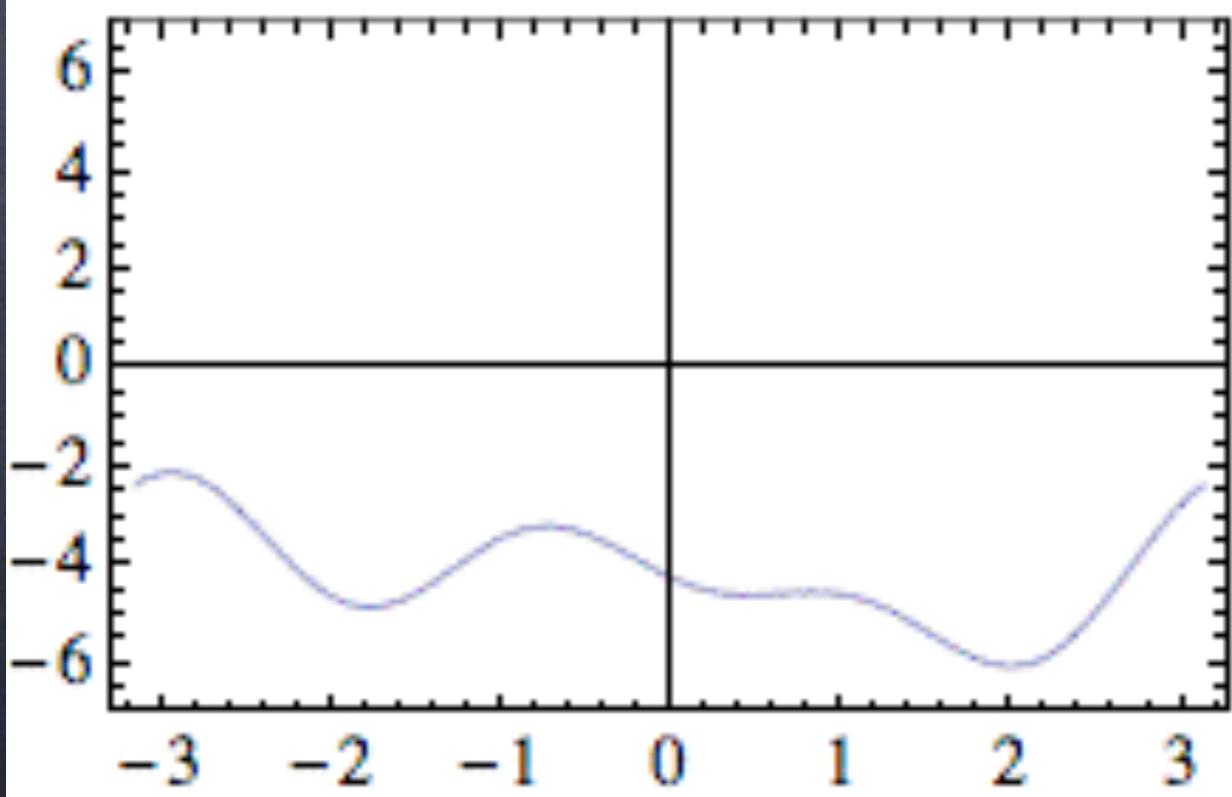
*note, with symmetric singularities at zero*

*adding a constant:*

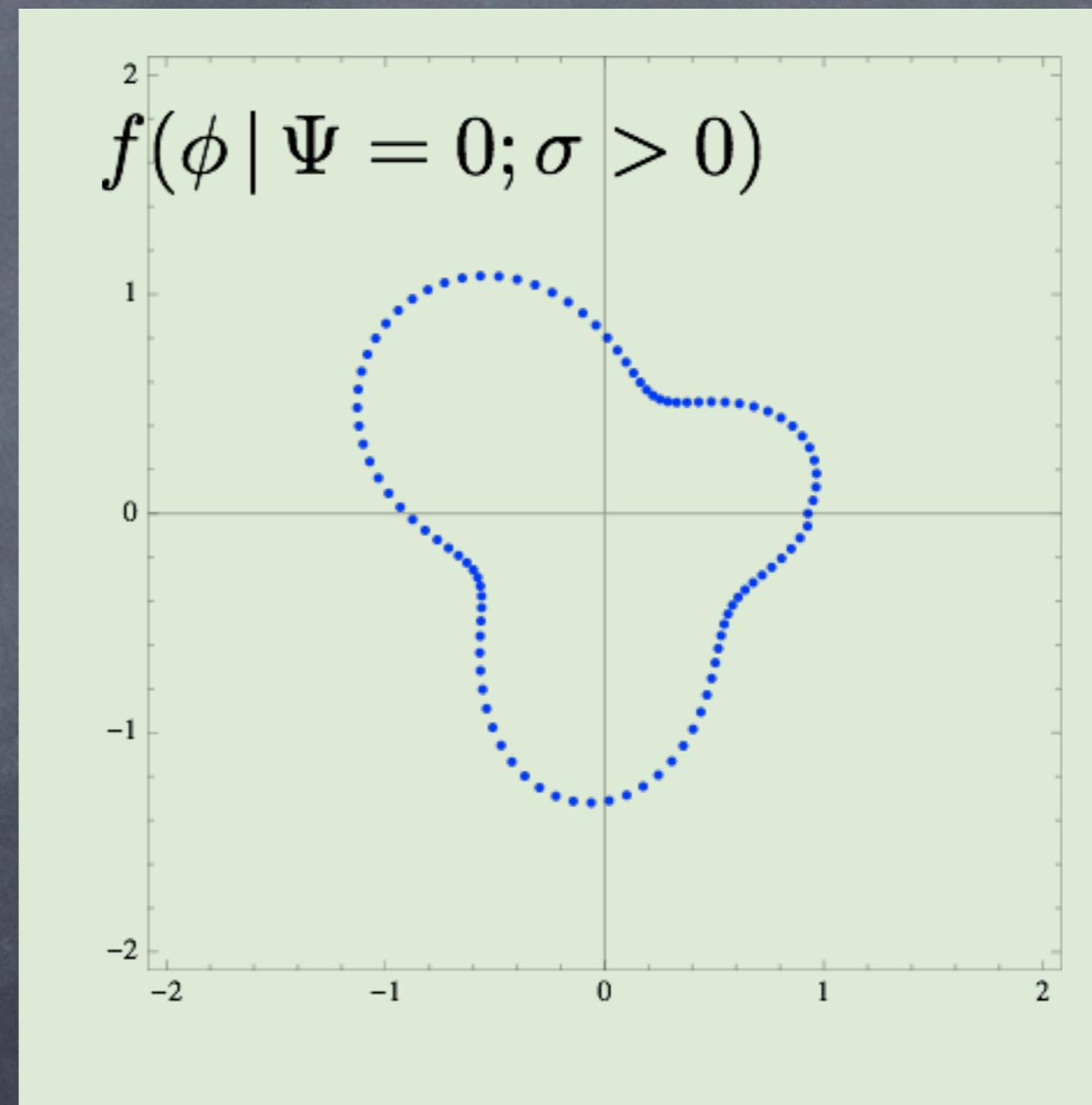
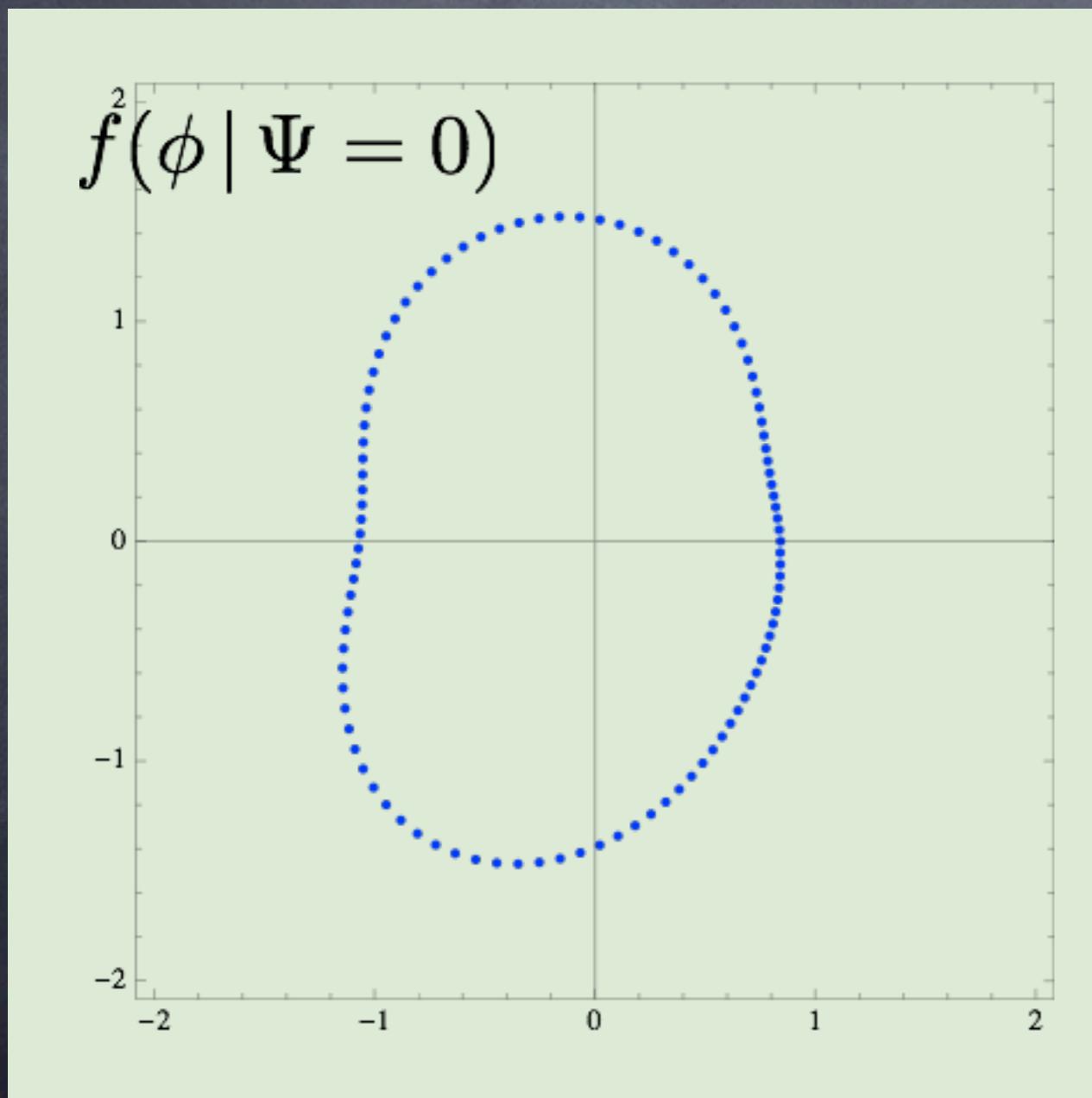
*adding a constant:*



*adding a constant:*



*a test run, cutting on handicity*



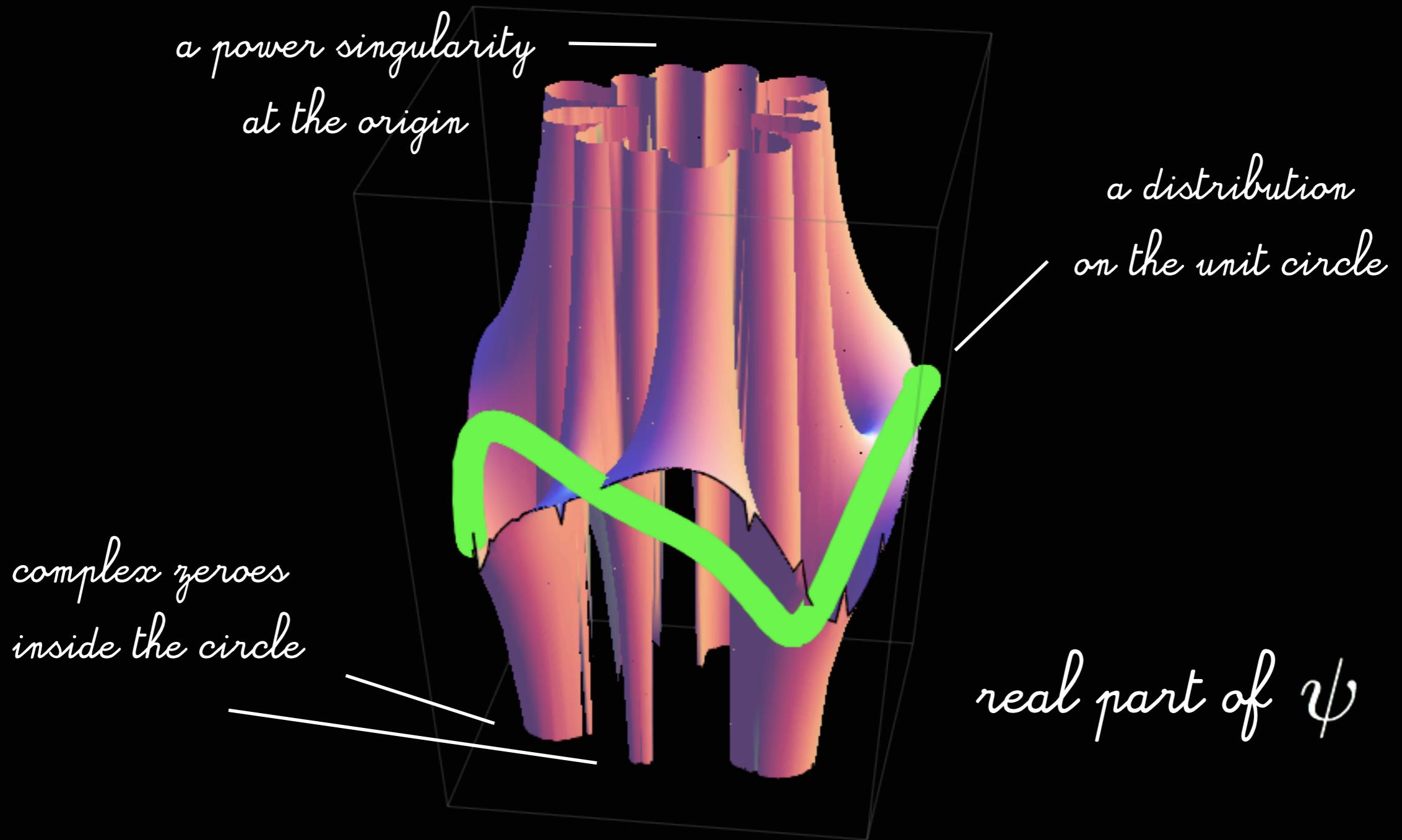


we propose to study strong interactions  
with event chirality sorting.

With transverse initial or inclusive  
final SPIN, this would be so  
**AWEOME**

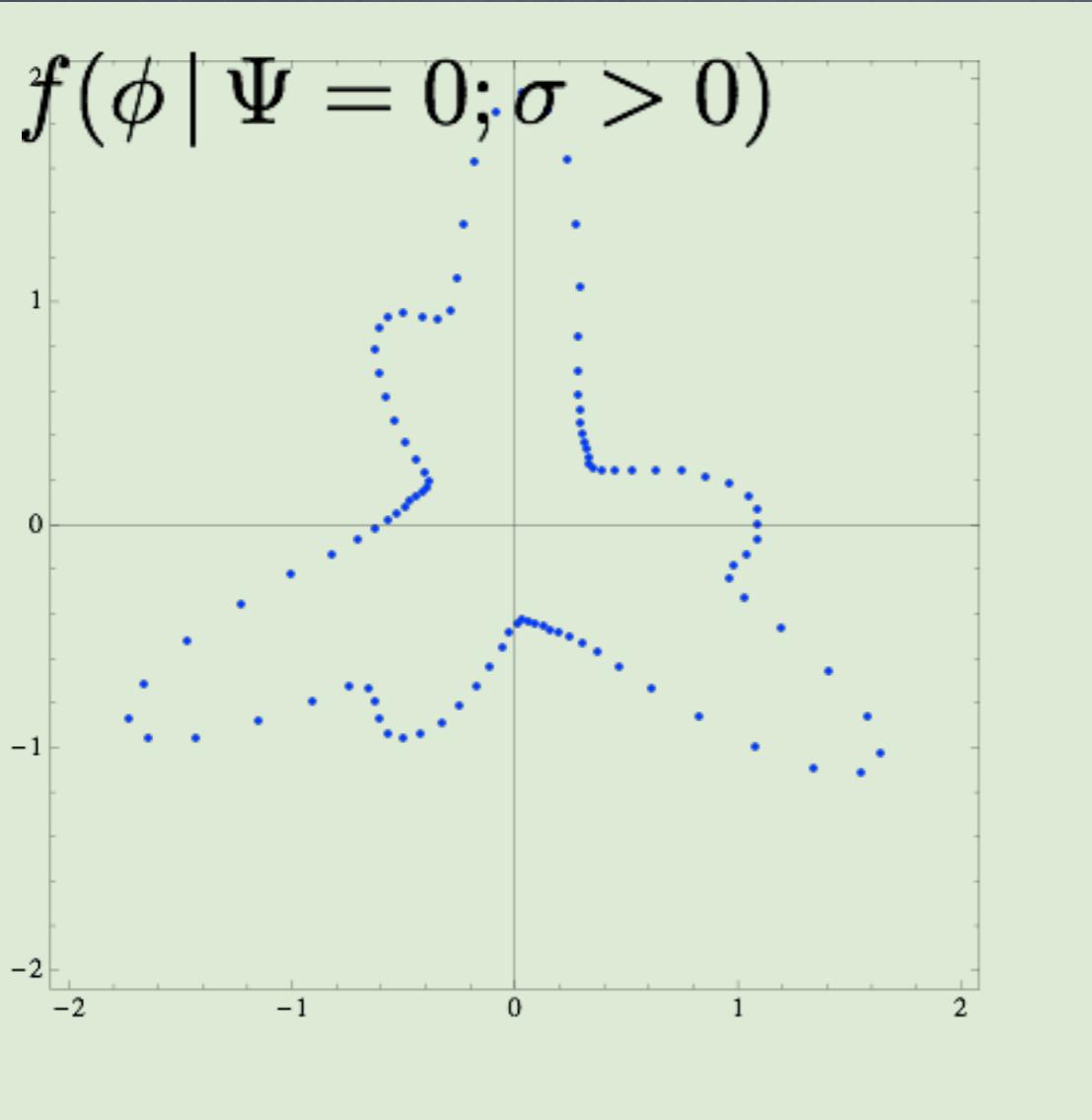
See our paper for high energy strong parity violation

What do those wave functions look like ?

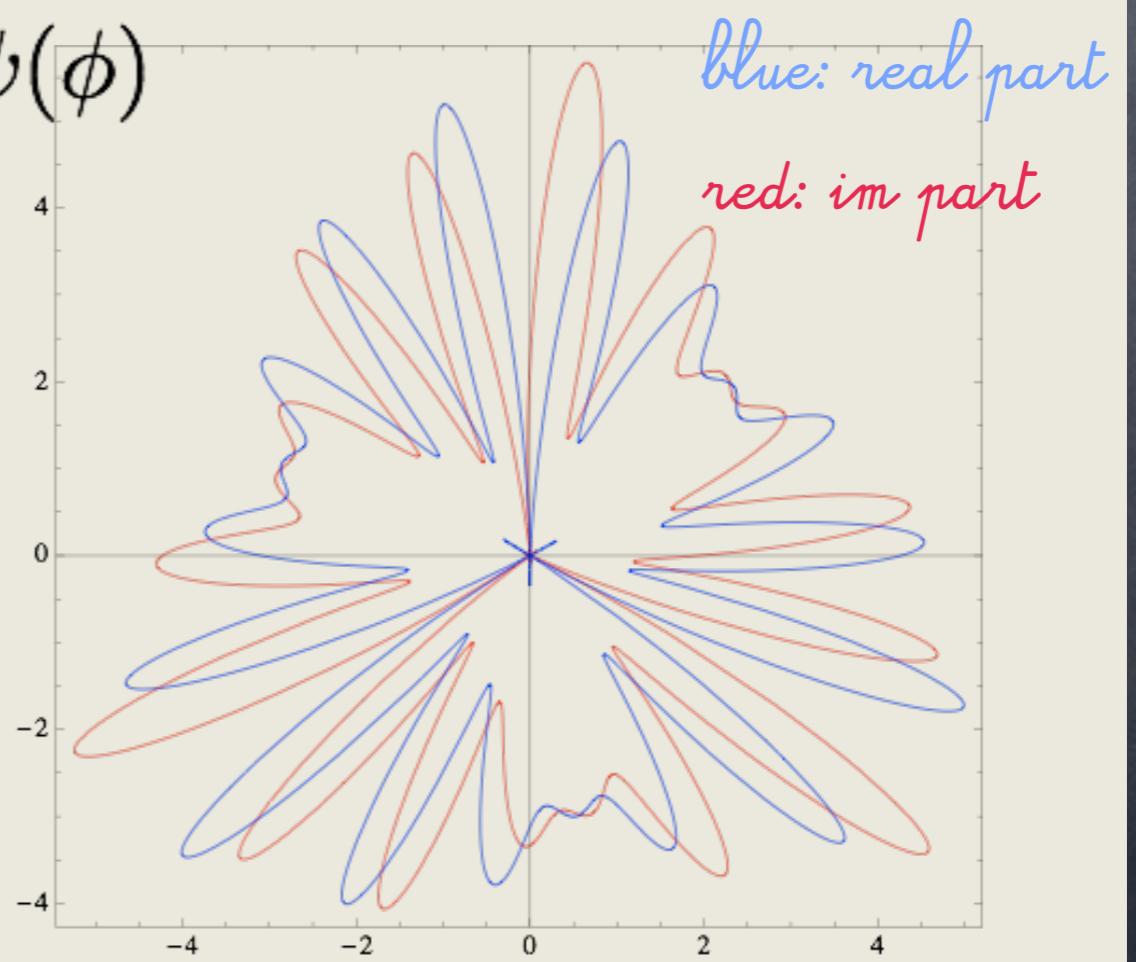


What do those wave functions look like ?

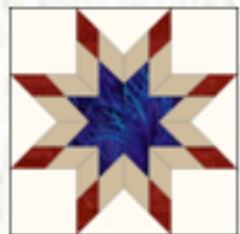
$$f(\phi \mid \Psi = 0; \sigma > 0)$$



$$\psi(\phi)$$

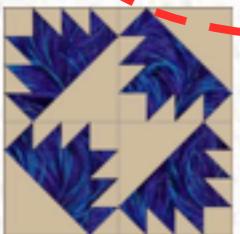


(8+, 8+)

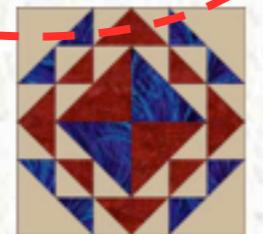


Blazing Star

### Complex Blocks

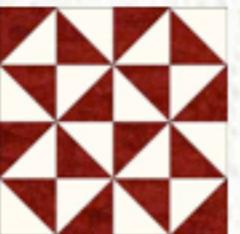


Barrister's Block



Corn and Beans Broken Dishes

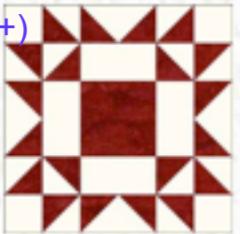
(4+, 4+)



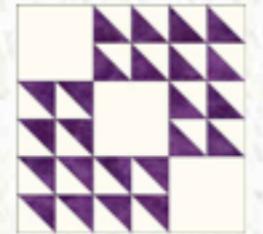
(4+, 4+)



Crown of Thorns



Arizona

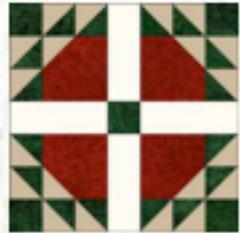


Cut Glass Dish

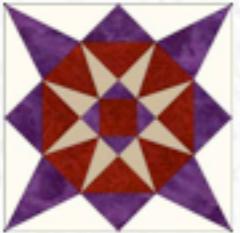


54-40 or Fight

(4+, 4-)



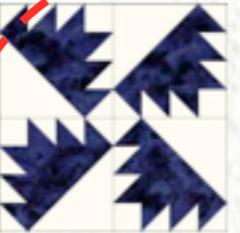
Dove in the Window Dogtooth Violet



Dogtooth Violet



Fool's Puzzle

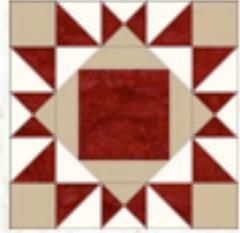


Kansas Troubles

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### Symmetry

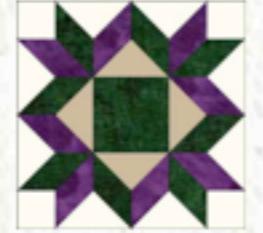
Hermann Weyl



King David's Crown



Lady of the Lake Laurel Wreath



Laurel Wreath

(1+, 0-+0+)

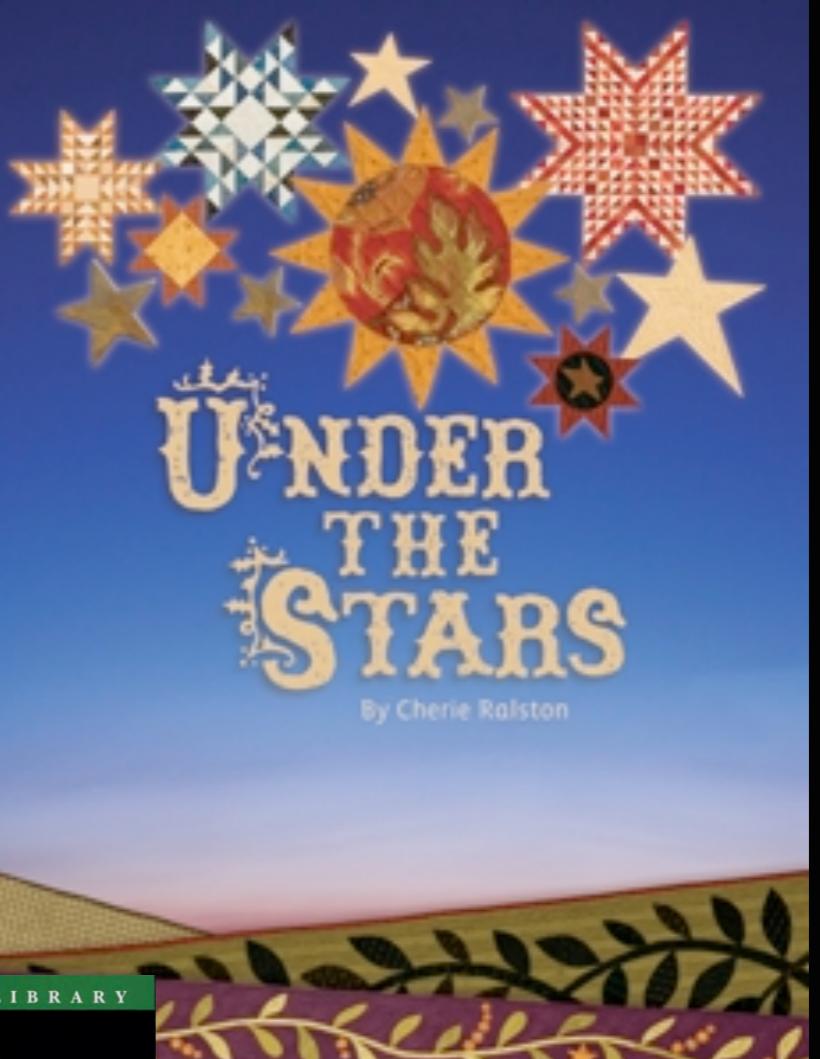
(4+, 4-)

(4+, 4-)

(4+, 4-)(4+, 4-)(4+, 4-)



Mariner's Compass

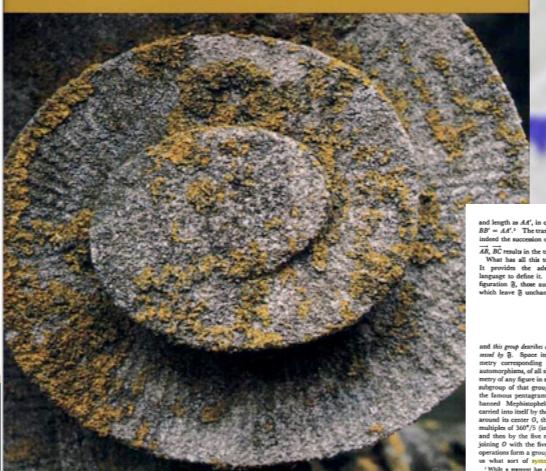


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and this group defines exactly the **symmetry group** of  $\Gamma$ . Specifically, it is the full group of all automorphisms of all orientations. The symmetry group of  $\Gamma$  is the group of all orientation preserving automorphisms of  $\Gamma$ , which is a subgroup of that group. Take for instance the square lattice  $\mathbb{Z}^2$ , which is often called the hatched Mephistopheles the devil. It is carried into itself by the five symmetries mentioned above, and also by all translations which are multiples of  $105^\circ$  (including the identity), and also by all rotations of  $105^\circ$  around the vertices. This ten symmetries form the **symmetry group** of  $\mathbb{Z}^2$ .  
What is this group? It is the group of all orientation preserving automorphisms of  $\mathbb{Z}^2$ , which is the group of all orientation preserving automorphisms of  $\mathbb{Z}^2$ .  
What has all this to do with **symmetry**?  
It goes back to the very beginning of mathematics, when we learned to count. We learned to count in the language to define it. Given a spatial configuration  $\mathcal{G}$ , those symmetries of  $\mathcal{G}$  which leave  $\mathcal{G}$  unchanged form a group  $\Gamma_{\mathcal{G}}$ .



FIG. 21

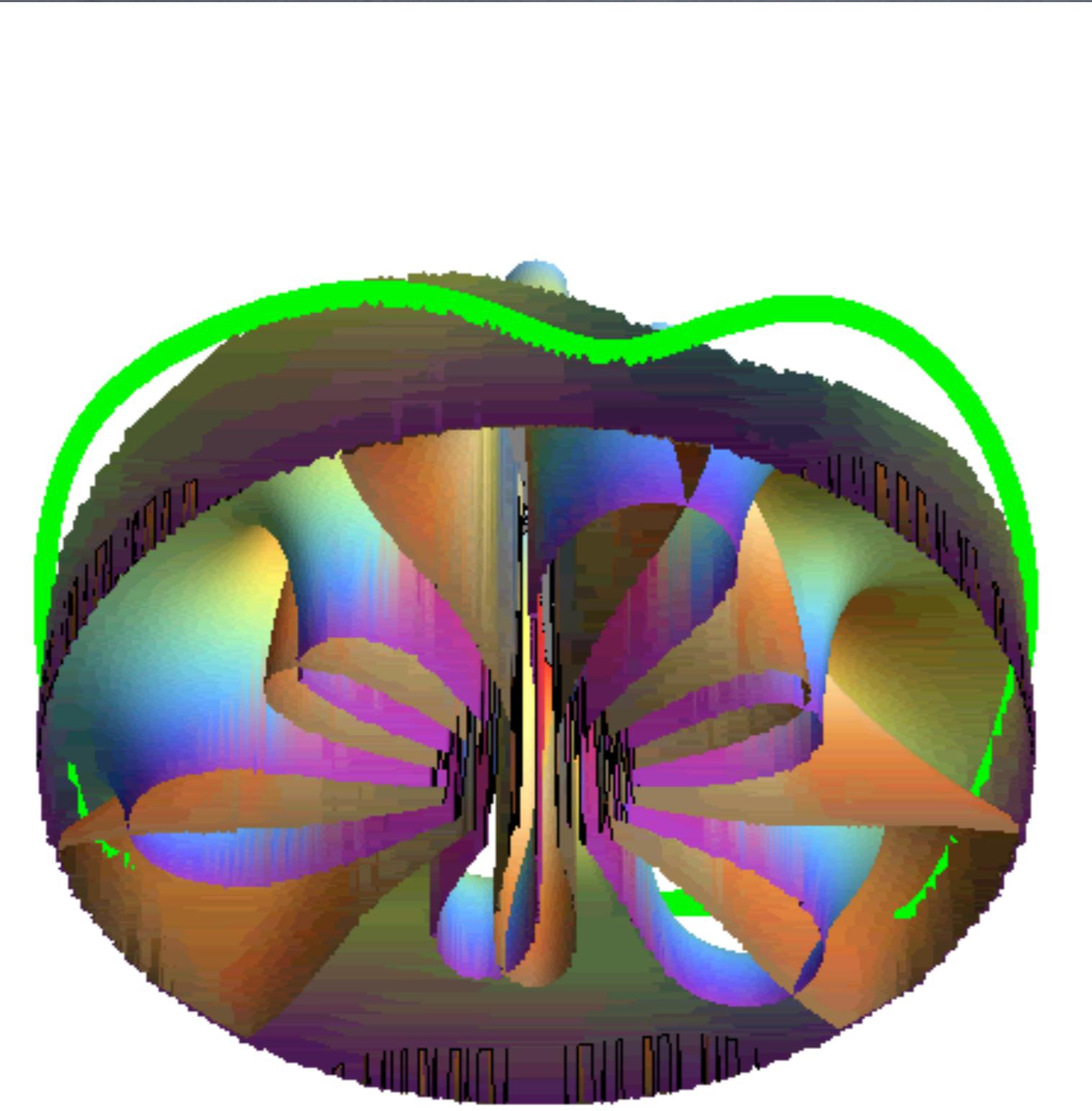


and length as  $A\theta'$ , in other words the vector  $AB = A\theta'$ . The translation forms a group; indeed the sum of two such translations  $AB, BC$  results in the translation  $AC$ .  
What has all this to do with **symmetry**?  
It goes back to the very beginning of mathematics, when we learned to count. We learned to count in the language to define it. Given a spatial configuration  $\mathcal{G}$ , those symmetries of  $\mathcal{G}$  which leave  $\mathcal{G}$  unchanged form a group  $\Gamma_{\mathcal{G}}$ .

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$D_4$  eigenstates  $(k, k')$ ;  $R_{4k=+}, P_{4k=-}$

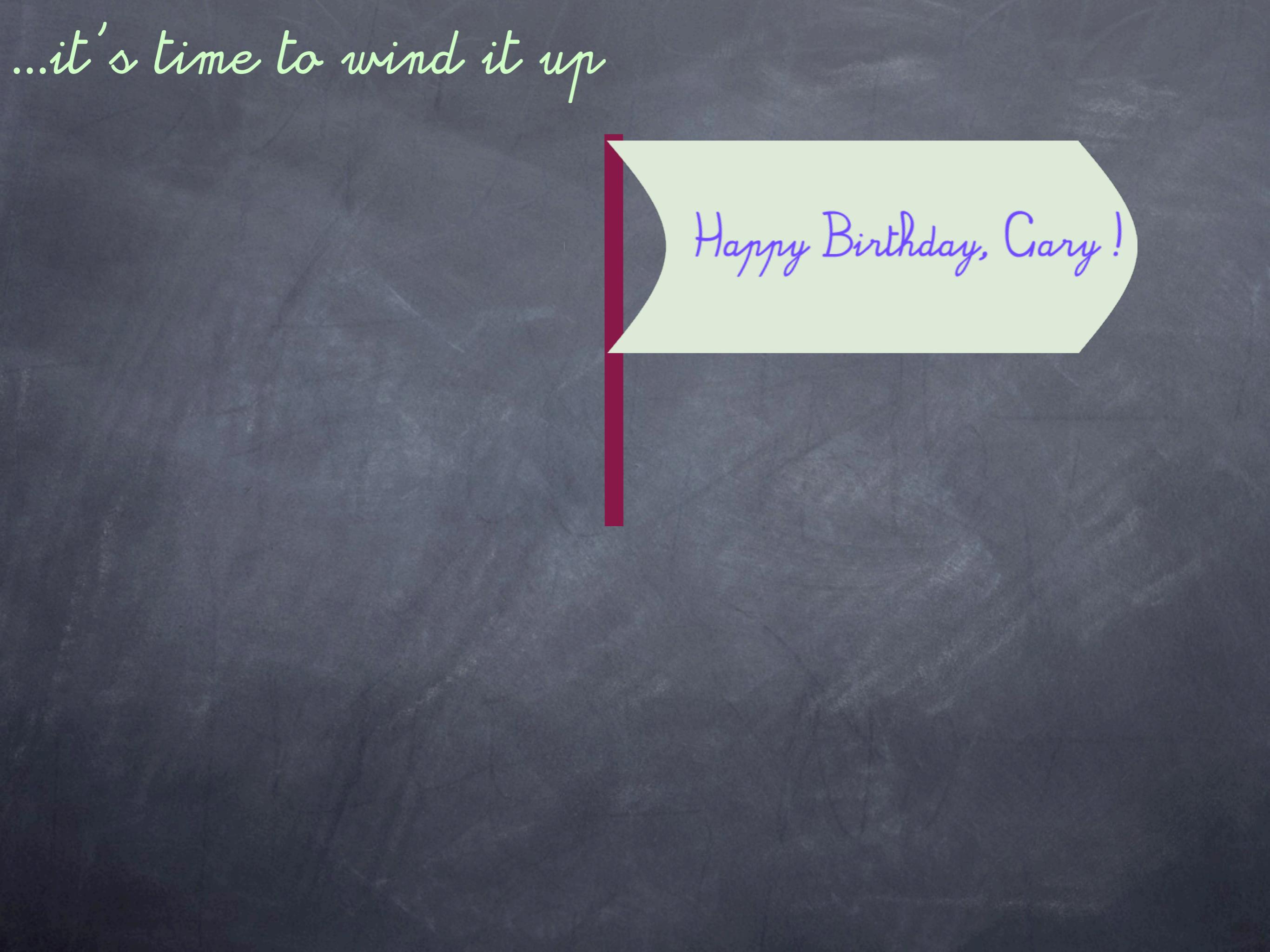
*...it's time to wind it up*



*...it's time to wind it up*



*...it's time to wind it up*



Happy Birthday, Gary !

*...it's time to wind it up*





Happy Birthday, Gary!